

Algebraic techniques for enumerating self-avoiding walks on the square lattice

This article has been downloaded from IOPscience. Please scroll down to see the full text article.

1993 J. Phys. A: Math. Gen. 26 1519

(<http://iopscience.iop.org/0305-4470/26/7/012>)

View [the table of contents for this issue](#), or go to the [journal homepage](#) for more

Download details:

IP Address: 171.66.16.68

The article was downloaded on 01/06/2010 at 21:04

Please note that [terms and conditions apply](#).

Algebraic techniques for enumerating self-avoiding walks on the square lattice

A R Conway†, I G Enting†§ and A J Guttmann‡||

† Department of Mathematics, The University of Melbourne, Parkville, Victoria 3052, Australia

‡ Department of Theoretical Physics, Oxford University, 1 Keble Road, Oxford, OX1 3NP, UK

Received 25 September 1992

Abstract. We describe a new algebraic technique for enumerating self-avoiding walks on the rectangular lattice. The computational complexity of enumerating walks of N steps is of order $3^{N/4}$ times a polynomial in N , and so the approach is greatly superior to direct counting techniques. We have enumerated walks of up to 39 steps. As a consequence, we are able to accurately estimate the critical point, critical exponent, and critical amplitude.

1. Introduction

Over the years, the enumeration of square lattice self-avoiding walks has become a benchmark, first for computer performance, and more recently for algorithm design. In the early 1970's, Sykes *et al* (1972) obtained 24 terms of the series by using the chain counting theorem. Direct enumeration is probably somewhat faster, but the graphs enumerated for the chain counting theorem (figure-eights, theta graphs, dumbbells and polygons) were useful for other problems in the theory of phase transitions, most notably the Ising model. Guttmann (1987) extended the series by three terms, using direct enumeration, and Guttmann and Wang (1991) using a dimerization algorithm obtained two further terms. Subsequently McDonald *et al* (1992) obtained a further term, using an extension of dimerization to trimerization. Late in 1991, Masand *et al* (1992) used a large supercomputer, a CM-2 containing 65 536 processors, to extend the series to 34 terms, running for ~ 100 hours. All these advances came about because of improvements in computer technology (in large part), with relatively small improvements brought about by algorithm design. (Dimerization or trimerization saves a factor of around 2 or 3, but does not ameliorate the exponential growth rate of computer time.)

The finite-lattice method plus transfer matrices described here allows 35 terms to be obtained on a work station (an IBM 6000/530 with 256MB of memory) in less time than the 65 536 processor CM-2 took to obtain 34 terms. Because of memory

§ Permanent address: CSIRO, Division of Atmospheric Research, Private Bag 1, Mordialloc, Victoria 3195, Australia.

|| Permanent address: Department of Mathematics, The University of Melbourne, Parkville, Victoria 3052, Australia.

requirements, it was necessary to move to a larger machine (an IBM 3090/400 with 500MB memory and 2GB of backing storage) to go from 35 to 39 steps. This computer would be capable of extending the series to 43 terms. However that calculation might take up to a month of CPU time, and so has not been pursued. This improvement however is due to the exponential improvement in algorithm design, rather than evolution of computer speed. This is discussed in more detail below.

The method used is based on the method which Enting and Guttmann have used extensively over the past twelve years to enumerate self-avoiding rings on a square lattice, but is significantly more complicated due to the requirement for a second stage of processing.

In a series of papers we have reported some significant improvements in the enumeration of self-avoiding rings on the square lattice extending the known series (Sykes *et al* 1972) from 26 steps to 38 steps (Enting 1980), 46 steps (Enting and Guttmann 1985) and 56 steps (Guttmann and Enting 1988). The extension from 38 steps to 56 steps reflects our use of increasingly powerful computing systems and in particular the use of increasingly large amounts of physical and virtual memory. Only minor changes to our programs have been made, primarily to 'tune' the procedure to make efficient use of particular computer architectures. We have also generalized our method to enumerate caliper moments of self-avoiding rings on the square lattice (Guttmann and Enting 1988) and also to enumerate self-avoiding rings on the L and Manhattan lattices (Enting and Guttmann 1985) and the honeycomb lattice (Enting and Guttmann 1989). Most recently we have extended the enumeration of triangular lattice polygons to 25 steps (as reported by Enting and Guttmann 1990) and then to 35 steps (Enting and Guttmann 1992).

While our techniques have been highly efficient for enumerating self-avoiding rings they are less suitable for enumerating self-avoiding walks. The difficulty is that walks can span a larger lattice than rings because they are not forced to return. A walk of L steps can span a distance L while a ring of L steps can only span a distance of up to $L/2$. For walks constrained by a surface our polygon enumeration techniques could be generalized in analogy with our calculation of surface susceptibilities for the square lattice (Enting and Guttmann 1980).

The present paper presents an algebraic technique for enumerating self-avoiding walks on a rectangular lattice. The basic quantity that we consider is $C_{m,n}$, the number of walks from a given origin with n steps in the $\pm x$ directions and m steps in the $\pm y$ directions. We consider segments of walks that double back in the y -direction and which can therefore be counted efficiently by transfer matrix techniques. The general enumeration can be expressed as a combination of such irreducible contributions. We further improve the efficiency of the procedure by restricting the range of the index n to $n \leq k$ and reconstruct $C_{m,n}$ for $m+n \leq 2k+1$ by using the symmetry relation

$$C_{m,n} = C_{n,m}. \quad (1.1)$$

The layout of the remainder of this paper is as follows. Section 2 describes the way in which the generating function for self-avoiding walks can be constructed from irreducible contributions. Section 3 shows the way in which these irreducible contributions can be constructed from generating functions for walks on strips that can be determined by algebraic techniques. Section 4 describes the algorithms for determining the requisite generating functions and analyses the computational complexity of the procedure. Section 5 describes our analysis of the singularity

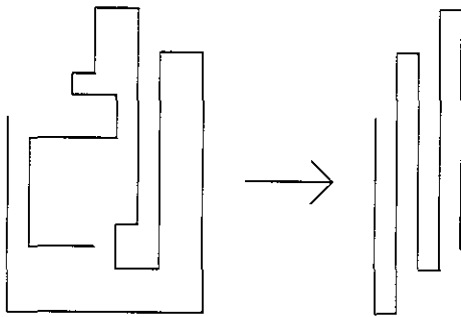


Figure 1. Schematic representation of projections of self-avoiding walks onto the y -axis.

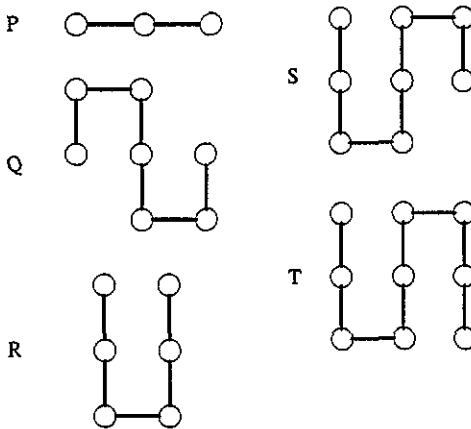


Figure 2. The five types of irreducible component from which self-avoiding walks are constructed.

structure of the generating function for self-avoiding walks based on the 39 terms that we have obtained.

2. Generating functions for self-avoiding walks

The generating function for self-avoiding walks on the rectangular lattice is

$$C(u, w) = \sum_{m, n=0}^{\infty} C_{mn} w^m u^n. \tag{2.1}$$

The enumeration of the coefficients C_{mn} is restricted to finite-order in m and/or n and we generally truncate the double series at $m + n \leq J$. The obvious summation then gives us the number of walks of up to J steps on the square lattice.

Our enumeration procedure is based on considering projections of walks onto the y -axis to produce the type of diagrams shown in figure 1. We refer to segments of walks as irreducible if the projection of that segment onto the y -axis has two or more y -bonds in each position. We will also classify irreducible segments by the number of y -bonds that they span. At this point we need to refine the terminology and distinguish between walks which are directed graphs and chains which are not directed. Our aim is to enumerate walks while the transfer matrix techniques generally enumerate chains. The decomposition of walks into irreducible components makes use of both chains and walks.

Figure 2 shows the five distinct types of irreducible component that we need to consider.

$P(u, w)$ is the generating function for walks that have no y -bonds, thus

$$P(u, w) = 1 + 2u + 2u^2 + 2u^3 + \dots = (1 + u)/(1 - u). \quad (2.2)$$

$Q(u, w)$ is the generating function for chains that are irreducible and for which neither end-point lies at an extremal y -coordinate. We also consider subdividing such cases according to m , the number of y -bonds spanned by the projection and define $Q_m(u, w)$ accordingly. The two pertinent results are

$$Q(u, w) = \sum_{m=2}^{\infty} Q_m(u, w) \quad (2.3)$$

and

$$Q_m(u, w) = O(w^{2m}). \quad (2.4)$$

Note that from the definition

$$Q_0 \equiv Q_1 \equiv 0. \quad (2.5)$$

$R(u, w)$ is the generating function for irreducible chains with both endpoints having the maximal y -coordinate. Again we subdivide these chains according to the number of y -bonds spanned and put

$$R(u, w) = \sum_{m=1}^{\infty} R_m(u, w) \quad (2.6)$$

$$R_m(u, w) = O(w^{2m}) \quad (2.7)$$

and arbitrarily define

$$R_0 \equiv 0. \quad (2.8)$$

$S(u, w)$ is the generating function for irreducible chains that have precisely one end having the maximal y -coordinate. Again

$$S(u, w) = \sum_{m=2}^{\infty} S_m(u, w) \quad (2.9)$$

$$S_m(u, w) = O(w^{2m+1}) \quad (2.10)$$

and from the definition

$$S_0 \equiv S_1 \equiv 0. \quad (2.11)$$

Finally $T(u, w)$ is the generating function for irreducible chains in which the two ends have maximal and minimal y -coordinates. We arbitrarily define

$$T_0 \equiv 0 \quad (2.12)$$

so that

$$T(u, w) = \sum_{m=1}^{\infty} T_m(u, w) \quad (2.13)$$

$$T_m(u, w) = O(w^{3m}). \quad (2.14)$$

When constructing the generating functions for walks we combine the irreducible components represented by P, Q, R, S and T by linking them with single y -bonds. These each contribute a factor of w to the generating function $C(u, w)$. Two single y -bonds that are adjacent in the projection are actually connected by a type P irreducible component. A chain whose projection is k consecutive single y -bonds will have a factor of P at each internal point. Note that the *walk* generating function P is required because each distinct direction along the x -axis will generate a distinct chain when combined with y -bonds. If we sum these contributions we can consider linking components of types R, S and T with chains of one or more y -bonds connecting type P walks. The generating function for such chains is

$$\begin{aligned} U(u, w) &= w + wPw + wPwPw + \dots \\ &= w/(1 - wP(u, w)) \end{aligned} \quad (2.15)$$

In the same way we can regard the overall chain as consisting of end segments P, R or S connected by combinations of 'reducible' parts with generating function U and irreducible parts with generating function T . The generating function for chains connecting irreducible end segments is thus

$$\begin{aligned} V(u, w) &= U + UTU + UTUTU + \dots \\ &= U/(1 - TU) \\ &= w/(1 - w(T + P)). \end{aligned} \quad (2.16)$$

It is now possible to express the self-avoiding walk generating function as

$$\begin{aligned} C(u, w) &= P(u, w) + 2[Q(u, w) + 2R(u, w) + 2S(u, w) + T(u, w)] \\ &\quad + 2V(u, w)[P(u, w) + 2R(u, w) + S(u, w) + T(u, w)]^2. \end{aligned} \quad (2.17)$$

The structure of expression (2.17) shows that to obtain an expansion in powers of w , it is necessary to obtain Q, R, S , and T to the requisite order. If the irreducible generating functions Q_m, R_m, S_m and T_m are known for $m \leq M$ then $C(u, w)$ will be correct to w^{2M+1} (the first incorrect term arising from the absences of Q_{M+1} and R_{M+1}). The use of the symmetry relations (1.1) will give $C_{m,n}$ for $m+n \leq 4M+3$.

3. Combining generating functions for chains on strips

The transfer matrix techniques described in the next section produce generating functions for sets of walks confined to strips whose y -coordinates are bounded. Subject to these constraints, all chains are counted, not merely irreducible components. Thus there is a need to relate unrestricted generating functions to the restricted generating functions for irreducible components. In our previous enumerations of self-avoiding rings only linear combinations of different classes of graph were involved and so the restricted generating functions were linear combinations of unrestricted generating functions. The present formalism is more complicated because nonlinear relations are involved. We begin by considering $T_M^*(u, w)$, the generating function for chains whose y -coordinates span M bonds and which have one end at each y -extremum, that is, T_M^* is the generating function for bridges. The sum over M is denoted $T^*(u, w)$. By considering the appropriate subset of terms from (2.17) we have

$$T^*(u, w) = T(u, w) + (P(u, w) + T(u, w))^2 V(u, w). \quad (3.1)$$

While this equation is formally correct, it is unsuitable for relating T^* to T because of the fact that while T_m is of order w^{3m} , T_m^* is of order w^m . Thus to obtain T correct to w^K would require the calculation of T_m^* for $m \leq K$. This difficulty is avoided by introducing an extra variable z whose power corresponds to the width of the segment under consideration. We refer to functions including z as 'extended generating functions'. We define the extended generating function for irreducible bridges as

$$X(u, w, z) = P(u, w) + \sum_{m=1}^{\infty} z^m T_m(u, w) \quad (3.2a)$$

and the extended generating function for all bridges as

$$X^*(u, w, z) = P(u, w) + \sum_{m=1}^{\infty} z^m T_m^*(u, w). \quad (3.2b)$$

In these terms, $V(u, w)$, the generating function for bridges ending in single bonds, generalizes to

$$\hat{V}(u, w, z) = X(u, w, z) = wz/[1 - wzX(u, w, z)]. \quad (3.3)$$

Relation (3.1) generalizes to

$$X^*(u, w, z) = X(u, w, z) + X(u, w, z)^2 \hat{V}(u, w, z) \quad (3.4)$$

whence

$$X(u, w, z) = X^*(u, w, z)/[1 + wzX^*(u, w, z)]. \quad (3.5)$$

This relation provides the basis of a suitable truncation. The expansion of (3.5) to order z^K requires T_m^* for $m = 1$ to K and will give T_m for $m = 1$ to K . The result, noted above, that T_m is of order w^{3m} provides a useful check on the algebra.

If we define $R_m^*(u, w)$ as the generating function for chains in a strip of width m such that both ends have the maximal y -coordinate then

$$R_m^*(u, w) = \frac{1}{2}(P(u, w) - 1) + \sum_{n=1}^m R_n(u, w). \tag{3.6}$$

This relation can be easily inverted to give individual R_m . These are needed to define $R(u, w)$ in (2.17) and also to recover the $S_m(u, w)$ from $Y(u, w, z)$ (3.8a) and the $Q_m(u, w)$ from $Z(u, w, z)$ (3.13).

We define $S_m^*(u, w)$ as the generating function for chains in a strip of width m where one end of the chain has the maximal y -coordinate and the other end does not have an extremal y -coordinate. The general relation between the irreducible and unrestricted generating functions is

$$S^*(u, w) = S(u, w) + [S(u, w) + 2R(u, w)]V(u, w)[P(u, w) + T(u, w)]. \tag{3.7}$$

We define the extended generating function for irreducible components with one or both ends at the maximal y co-ordinate as

$$Y(u, w, z) = 2 \sum_{m=1}^{\infty} z^m R_m(u, w) + \sum_{m=2}^{\infty} z^m S_m(u, w) \tag{3.8a}$$

and the corresponding unrestricted function as

$$Y^*(u, w, z) = 2 \sum_{m=1}^{\infty} z^m R_m^*(u, w) + \sum_{m=2}^{\infty} z^m S_m^*(u, w) \tag{3.8b}$$

Equation (3.7) generalizes to

$$Y^*(u, w, z) = Y(u, w, z)[1 + \tilde{V}(u, w, z)X(u, w, z)] \tag{3.9}$$

or

$$Y(u, w, z) = Y^*(u, w, z)/[1 + \tilde{V}(u, w, z)X(u, w, z)] \tag{3.10}$$

from which the S_m can be recovered once R_n and T_n are known for $n \leq m$.

Finally we consider Q_m^* , the generating function for chains in a strip X where neither end has an extremal y -coordinate.

We have

$$Q^*(u, w) = Q(u, w) + [P(u, w) + 2R(u, w)]^2V(u, w) \tag{3.11}$$

defining

$$Z(u, w, z) = \sum_{m=2}^{\infty} z^m Q_m(u, w) \tag{3.12a}$$

and

$$Z^*(u, w, z) = \sum_{m=2}^{\infty} z^m Q_m^*(u, w) \tag{3.12b}$$

gives

$$Z(u, w, z) = Z^*(u, w, z) - Y(u, w, z)^2\tilde{V}(u, w, z). \tag{3.13}$$

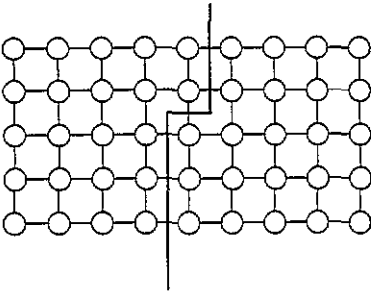


Figure 3. The cross section line defining the set of lattice bonds which specify the partial generating functions.

4. Transfer matrix enumeration techniques

The analysis in the previous sections has reduced the problem of enumerating general self-avoiding walks of $4K + 3$ steps to one of enumerating walks confined to strips of width $\leq K$ subject to various constraints. In order to enumerate walks confined to strips we use a transfer matrix technique that generalizes the approach that we have used in our earlier enumeration of self-avoiding rings. We draw a cross section line (with a kink) across the width of a strip of width K so as to cut $K + 2$ of the bonds on which steps of chains can occur (figure 3). We note that if we specify the set of occupied steps then the self-avoidance constraint acts independently to the left and right of the cross section line. However not all combinations of self-avoiding components from the left and right of the cross section line combine to give walks. It is necessary to consider the connectivity of the components. This can be done by generalizing the technique that we used in our earlier work. We assign to each bond intersected by the cross section line an index

$$n_i = 0, 1, 2 \text{ or } 3 \quad i = 1 \text{ to } K + 2.$$

Here '0' denotes an empty bond, '1' denotes a step connected to a (uniquely defined) later step, '2' denotes a step connected to a (uniquely defined) earlier step and '3' denotes a step not connected to any other steps intersected by the cross section line.

If we define

$$A(i, j) = \{k : k \leq j \text{ and } n_k = i\} \quad (4.1)$$

then we require

$$|A(1, j)| \geq |A(2, j)| \quad \text{for all } j \quad (4.2a)$$

$$|A(1, K + 2)| = |A(2, K + 2)| \quad (4.2b)$$

as in the enumeration of self-avoiding rings and

$$|A(3, K + 2)| \leq 2. \quad (4.2c)$$

The numbers of sets of n_i subject to these constraints for various K are given in table 1. These numbers give the main limitation on the size of walks that can be obtained because it is necessary to store a partial generating function for the number of walks corresponding to each set of n_i allowed. These numbers, s_k , are larger than

Table 1. The sizes of vectors required by the transfer matrix formalism. For ring enumeration r_k components are required. For walk enumeration s_k components are required.

Strip width k	#bonds	r_k	s_k
—	1	1	2
—	2	2	5
1	3	4	13
2	4	9	37
3	5	21	106
4	6	51	312
5	7	127	925
6	8	323	2 767
7	9	835	8 314
8	10	2 188	25 073
9	11	5 798	75 791
10	12	41 835	229 495

Table 2. Allowed transformations of bond indices. For the operations, 'Build' means incorporate the contribution into the new vector, as defined by the new indices. 'R($a \rightarrow b$)' means apply the change $a \rightarrow b$ to the other end of the chain in order to specify the index in the new vector. 'Ignore' means perform no operation, as a disconnected ring has been generated. 'Acc' means accumulate the vector component into the chain generating function if all the other n_i are zero, otherwise the operation preceding the 'OR' is applied.

Old indices (n_j, n_{j+1})	New indices (n_j, n_{j+1})	Operation
(0,0)	(0,0), (0,3), (3,0) or (1,2)	Build
(0,1) or (1,0)	(0,1), (1,0) or (0,0)	R(2 \rightarrow 3)
(0,2) or (2,0)	(0,2), (2,0) or (0,0)	R(1 \rightarrow 3)
(0,3) or (3,0)	(0,3), (3,0)	Build OR Acc
(1,1)	(0,0)	R(2 \rightarrow 1)
(2,2)	(0,0)	R(1 \rightarrow 2)
(1,2)	—	Ignore
(2,1)	(0,0)	
(3,3)	—	Ignore OR Acc
(3,1), (1,3)	(0,0)	R(2 \rightarrow 3)
(3,2), (2,3)	(0,0)	R(1 \rightarrow 3)

the corresponding vector sizes, r_k , used in the enumeration of self-avoiding walks, but only by a factor $\gamma(k)$ which is constrained as

$$1 \leq \gamma(k) \leq \frac{1}{2}(k^2 + 5k + 7).$$

Thus the increase in the s_k is dominated by a $3^{k/4}$ increase as for the r_k .

Self-avoiding chains in strips are developed successively by advancing the cross section line so that one vertex of the lattice passes from the right to the left of the line. Except at the beginning of a column, this corresponds to moving the kink down one row. This move replaces two bonds (and their associated n_i) by two new bonds with new n_i . The other n_i are unchanged except when the addition of the new site changes the connectivity of the components. Table 2 shows the various combinations of new n_i that can be produced from various combinations of old n_i .

pairs. Various special cases occur. If a link from a free end (i.e. $n_i = 3$) connects to an existing loop segment then the other end of the loop must be reset to type 3. If two loops meet then one end must be relabelled. The closing of a single loop implies that a ring, disconnected from the walk, has been created and this configuration is ignored. The final step in the construction of a chain is when two type 3 bonds meet. When this occurs, all other bonds must be empty for a valid chain generating function to be added to the running total. If this is not the case it implies that other disconnected components are present and the configuration is ignored.

The iteration is initiated from an empty state ($n_i \equiv 0$) with generating function 1. As each bond is added, factors of v or w as appropriate are used to multiply the old partial generating function and the product is accumulated into the running total for the new partial generating function. Each chain could, in principle, be generated for a number of different x - and y -displacements from a given reference origin. To ensure uniqueness in the x -direction, the chain is required to intersect the first column considered. Thus the state with all n_i zero is never continued after the first column has been built up. This ensures that each chain is counted only once and, together with the requirement that the last operation is to join two type 3 bonds, also ensures that only one connected component occurs in each graph that is counted. Formally, if $4k + 3$ series terms are required then the transfer matrix operation must be repeated until $4k + 3$ columns of each strip have been generated. This will ensure that all the cancellations involved in going from unrestricted to irreducible contributions will be correct. As noted above, the cancellations provide a useful check on the implementation of the algebraic formalism. If however the check is not required then the results (2.4), (2.7), (2.10) and (2.14) can be assumed to be true. For a strip of width M it is sufficient to generate only $4k + 3 - 2M$ columns of the strip.

This requires a large amount of memory to store all the intermediate generating functions. If this amount of memory is not available as physical memory, but only as virtual (disk based) memory, the performance can be enhanced greatly by being careful of the order in which the partial generating functions are processed. This can make the entire process close to sequential, and enormously reduce the amount of disk access required.

As mentioned before, when a partial generating function is processed, only the two n_i coefficients adjacent to the site being added will change, unless this changes the connectivity of a loop. However in this case, it will always be a non-zero going to another non-zero, so if we make up a 'partial signature' out of the n_i coefficients that are not being directly changed, and just record whether they are zero or not, this partial signature will be invariant under the transfer matrix. That is, all 'new' partial generating functions will have the same 'partial signature'.

We can then process all the partial generating functions with one particular partial signature, save the results to disk, and process the next set, until all are processed. We do not need to worry about the possibility of having to accumulate a new partial generating function to one stored on disk, because if two partial generating functions have different partial signatures, they definitely cannot have the same 'full signature' (n_i values).

The only problem with this is to make sure that we process all the partial generating functions of a particular partial signature first. Fortunately, one can get away without having to sort the output file before using it as input. If instead of saving to just one output file, two output files are used, it is possible to arrange

things such that one only has to read from these two files in order, and the data will be ordered in exactly the right way for adding the next site. This minimizes the amount of disk access. To obtain this nice ordering, one looks at whether the new value of n_i which will be included in the partial signature for the next phase is zero or not, and accordingly assigns it to one of the two output files. This will arrange the partial signatures in binary order, with the most significant bit in the partial signature being the most recent bit calculated, and the least significant bit being the oldest bit calculated.

This method has a side benefit—in a multi-processor architecture, one can have each processor working independently on separate partial signature groups, with very little inter-processor communication. This means that this algorithm is easy to parallelize, which will become increasingly important in the future as parallel computers are becoming more and more popular, whilst many other algorithms are difficult to parallelize.

The problem of ensuring uniqueness in the y -direction requires inclusion-exclusion arguments of the type used in our enumeration of self-avoiding rings. For a strip of width M we classify chain generating functions as $G_m(\pm, \pm, \pm)$ where $+$ denotes an allowed location for ends and $-$ denotes a forbidden location. The three arguments refer to the top row, the bottom row and the set of internal rows respectively. Specifying the end locations permissively, rather than prescribing how many end points lie in each set allows us to treat the two ends of the chain independently. We have

$$G_K(+, -, -) = \frac{1}{2}(P - 1) + \sum_{m=1}^K R_m \quad (4.3a)$$

$$G_K(+, +, -) = T_K^* + P - 1 + 2 \sum_{m=1}^K R_m \quad (4.3b)$$

$$G_K(+, -, +) = \frac{1}{2}K(P - 1) + \sum_{m=1}^K (K + 1 - m)(R_m + Q_m^* + S_m^*) \\ + \sum_{m=1}^K (K - m)(T_m^* + R_m + S_m^*) \quad (4.3c)$$

$$G_K(-, -, +) = \frac{1}{2}(K - 1)(P - 1) + \sum_{m=1}^K (K + 1 - m)Q_m^* + 2 \sum_{m=1}^{K-1} (K - m)(R_m + S_m^*) \\ + \sum_{m=1}^{K-2} (K - 1 - m)T_m^* \quad (4.3d)$$

These relations can be explicitly inverted to give

$$R_m = G_m(+, -, -) - G_{m-1}(+, -, -) \quad (4.4a)$$

$$Q_m^* = G_m(-, -, +) - G_{m-1}(+, -, +) - \sum_{n=1}^{m-1} (Q_n^* + R_n + S_n^*) \quad (4.4b)$$

$$S_m^* = G_m(+, -, +) - G_{m-1}(+, -, +) - G_m(+, -, -) - G_{m-1}(+, -, -) - \frac{1}{2}(P - 1) - Q_m^* - \sum_{n=1}^{m-1} (Q_n^* + 2 * S_n^* + T_m^*) \tag{4.4c}$$

$$T_m^* = G_m(+, +, -) - 2G_m(+, -, -). \tag{4.4d}$$

5. Analysis of series

This algorithm was implemented for computers in two parts. The first part was the transfer matrix portion; that is it computed the *G* functions. This was the time and memory intensive section. It was run up to a width of eight on an IBM RS6000/530 with 256MB of RAM, then to width nine on an IBM 3090 with 500MB RAM and 2GB backing store. It required about 200MB and 600MB of memory respectively, and required several days in each case. The same machine could have been used to do width 10 (43 terms) given several weeks of processor time. This program was written in C, and reused memory whenever possible (there was only one bank of memory, used for both the about-to-be-processed partial generating functions, and the have-just-been-processed partial generating functions). The method described previously for using disk storage efficiently was implemented and tested, but not used as whilst memory was then no longer a problem, time became a large problem.

The second part performed all the algebra. Whilst algebra on large (over 15 000 coefficients) polynomials in three variables is slow, it is still a minor problem compared to the transfer matrix section, requiring time in only hours and memory in sub-megabytes, so efficiency was not as vital. It was implemented in C++.

The results for width 9 (39 terms) are given in table 3.

Table 3. Numbers of self-avoiding walks.

<i>n</i>	<i>c_n</i>	<i>n</i>	<i>c_n</i>	<i>n</i>	<i>c_n</i>
0	1	14	2 374 444	28	2 351 378 582 244
1	4	15	6 416 596	29	6 279 396 229 332
2	12	16	17 245 332	30	16 741 957 935 348
3	36	17	46 466 676	31	44 673 816 630 956
4	100	18	124 658 732	32	119 034 997 913 020
5	284	19	335 116 620	33	317 406 598 267 076
6	780	20	897 697 164	34	845 279 074 648 708
7	2 172	21	2 408 806 028	35	2 252 534 077 759 844
8	5 916	22	6 444 560 484	36	5 995 740 499 124 412
9	16 268	23	17 266 613 812	37	15 968 852 281 708 724
10	44 100	24	46 146 397 316	38	42 486 750 758 210 044
11	120 292	25	123 481 354 908	39	113 101 676 587 853 932
12	324 932	26	329 712 786 220		
13	881 500	27	881 317 491 628		

The method of analysis used is based on first- and second-order differential approximants. It was also used in previous papers by Guttmann (1987), Guttmann and Wang (1989) and is described in detail in Guttmann (1989). In summary, we construct

near-diagonal inhomogeneous approximants, with the degree of the inhomogeneous polynomial increasing from 1 to 8 in steps of 1. For first-order approximants ($K = 1$), 12 approximants are constructed that utilize a given number of series coefficients, N . Rejecting occasional defective approximants, we form the mean of the estimates of the critical point and critical exponent for fixed order of the series, N . The error is assumed to be two standard deviations. A simple statistical procedure combines the estimates for different values of N by weighting them according to the error, with the estimate with the smallest error having the greatest weight. As the error tends to decrease with the number of terms used in the approximant, this procedure effectively weights approximants derived from a larger number of terms more heavily.

Table 4. Estimates of the critical point (x_c) and critical exponent (γ) from first-order ($K = 1$) and second-order ($K = 2$) differential approximants. L is the number of approximants used. If L is too small (marked with an 'x'), the estimates are not used in the subsequent statistical analysis.

K	n	x_c	Error	γ	Error	L
1	19	0.379 0473	—	-1.343 0184	—	1x
1	20	0.379 0495	0.000 0004	-1.343 2502	0.000 1073	2x
1	21	0.379 0526	0.000 0105	-1.343 5231	0.001 6130	4
1	22	0.379 0469	0.000 0133	-1.342 7806	0.002 1290	4
1	23	0.379 0468	0.000 0071	-1.342 7666	0.001 1925	5
1	24	0.379 0520	0.000 0038	-1.343 6373	0.000 6432	7
1	25	0.379 0525	0.000 0052	-1.343 7004	0.001 0102	9
1	26	0.379 0508	0.000 0050	-1.343 3692	0.001 0094	11
1	27	0.379 0530	0.000 0074	-1.343 7617	0.001 3817	12
1	28	0.379 0519	0.000 0016	-1.343 5724	0.000 3713	12
1	29	0.379 0519	0.000 0010	-1.343 5632	0.000 2717	9
1	30	0.379 0517	0.000 0016	-1.343 5091	0.000 4340	9
1	31	0.379 0518	0.000 0011	-1.343 5230	0.000 3276	8
1	32	0.379 0514	0.000 0017	-1.343 4149	0.000 5073	4
1	33	0.379 0521	0.000 0016	-1.343 6098	0.000 4159	9
1	34	0.379 0525	0.000 0028	-1.343 7270	0.000 7447	11
1	35	0.379 0518	0.000 0003	-1.343 5417	0.000 0956	10
1	36	0.379 0519	0.000 0003	-1.343 5722	0.000 1233	10
1	37	0.379 0517	0.000 0011	-1.343 4730	0.000 4568	8
1	38	0.379 0518	0.000 0009	-1.343 5048	0.000 3501	9
1	39	0.379 0521	0.000 0001	-1.343 6392	0.000 0515	2x
2	28	0.379 0525	—	-1.343 7307	—	1x
2	29	0.379 0520	—	-1.343 5885	—	1x
2	30	0.379 0518	0.000 0006	-1.343 5431	0.000 1628	2x
2	31	0.379 0518	0.000 0005	-1.343 5300	0.000 1315	3x
2	32	0.379 0513	0.000 0016	-1.343 2041	0.001 0873	3x
2	33	0.379 0515	0.000 0007	-1.343 3885	0.000 4434	4
2	34	0.379 0519	0.000 0003	-1.343 5503	0.000 0824	5
2	35	0.379 0521	0.000 0004	-1.343 6142	0.000 1314	5
2	36	0.379 0519	0.000 0006	-1.343 5607	0.000 1972	6
2	37	0.379 0520	0.000 0001	-1.343 5822	0.000 0276	7
2	38	0.379 0520	0.000 0001	-1.343 5845	0.000 0264	6
2	39	0.379 0521	0.000 0002	-1.343 6174	0.000 0616	4

For second-order approximants ($K = 2$), we construct eight distinct approximants for each value of N . A summary of the results of this process is shown in table 4.

The statistical procedure used to combine the results gives

$$x_c = 0.379\,052 \pm 0.000\,001 \quad \gamma = 1.3435 \pm 0.0003 \quad (K = 1)$$

$$x_c = 0.379\,0520 \pm 0.000\,0005 \quad \gamma = 1.3436 \pm 0.000\,15 \quad (K = 2).$$

These results provide abundant support, if support is still needed, for the value $\gamma = 1.343\,75$ obtained by Nienhuis (1982, 1984). To refine the estimate of the critical point, linear regression is used. There is a strong correlation between estimates of the critical point and critical exponent. This is quantified by linear regression, and in this way the biased estimates (biased at $\gamma = 43/32$) are obtained.

We find

$$x_c = 0.379\,0524 \pm 0.000\,0005 \quad (K = 1)$$

$$x_c = 0.379\,0525 \pm 0.000\,0005 \quad (K = 2).$$

These are in excellent agreement with previous estimates based on the 56 term polygon series (Guttmann and Enting 1988), $x_c = 0.379\,052\,28 \pm 0.000\,000\,14$.

For the honeycomb lattice, the 'connective constant' $= 1/x_c$ is known exactly (Nienhuis 1982, 1984), and is $\sqrt{2 + \sqrt{2}}$, which satisfies a simple quadratic equation in x_c^2 . A feature of Maple (Version 5) is a clever algorithm for seeking polynomials with integer coefficients that have a given root. Attempting to find a quartic polynomial that gave as a root the biased value of x_c quoted above, we found the best solution was also a polynomial quadratic in x_c^2 . It was

$$581x^4 + 7x^2 - 13 = 0.$$

The root is $x_c = 0.37905227\dots$. While we consider it would be fortuitous if this were the true value of the critical point, it nevertheless provides a useful mnemonic.

Another analysis we were able to carry out with this long series was a study of amplitudes of the leading term and the correction terms. As previously discussed for self-avoiding polygons (Guttmann and Enting 1988), we have found no evidence for any non-analytic correction-to-scaling term other than that suggested by Nienhuis, with a 'correction' exponent of $\Delta = 1.5$. In the case of the polygon generating function this 'folds into' the additive analytic term. However, for the SAW series, it gives rise to a non-analytic correction term. Furthermore, there is another singularity on the negative real axis, at $x = -x_c$, as shown by Guttmann and Whittington (1978).

Thus we expect the generating function for walks to behave like

$$C(x) = \sum c_n x^n \sim A(x)(1 - \mu x)^{-43/32} [1 + B(x)(1 - \mu x)^{3/2} + \dots] \\ + D(x)(1 + \mu x)^{+1/2}. \quad (5.1)$$

The exponent for the singularity on the negative real axis reflects the fact that this term is expected to behave as the energy, and hence to have exponent $1 - \alpha$, where $\alpha = \frac{1}{2}$. From the above, it follows that the asymptotic form of the coefficients, c_n , behaves like

$$c_n \sim \mu^n [a_1 n^{11/32} + a_2 n^{-21/32} + b_1 n^{-37/32} + (-1)^n d_1 n^{-3/2} + (-1)^n d_2 n^{-5/2}]. \quad (5.2)$$

Table 5. Sequences of amplitude estimates. Refer to equation (5.2).

n	d_2	d_1	b_1	a_2	a_1
29	0.0639	-0.1878	-0.1999	0.5584	1.177 00
30	0.0666	-0.1879	-0.2022	0.5590	1.176 99
31	0.0715	-0.1881	-0.1980	0.5579	1.177 00
32	0.0738	-0.1882	-0.1999	0.5584	1.177 00
33	0.0781	-0.1883	-0.1963	0.5574	1.177 01
34	0.0800	-0.1884	-0.1979	0.5578	1.777 00
35	0.0838	-0.1885	-0.1947	0.5570	1.777 01
36	0.0855	-0.1885	-0.1960	0.5573	1.777 01
37	0.0890	-0.1886	-0.1932	0.5566	1.777 01
38	0.0904	-0.1887	-0.1943	0.5569	1.777 01
39	0.0936	-0.1888	-0.1919	0.5563	1.777 02

The five amplitudes, a_1, a_2, b_1, d_1 and d_2 come from the leading singularity (giving rise to a_1 and a_2), the correction-to-scaling term (giving rise to b_1) and the term on the negative real axis (giving rise to d_1 and d_2). A small program written in Mathematica was used to fit successive quintuples of coefficients, $c_{n-4}, c_{n-3}, c_{n-2}, c_{n-1}$ and c_n for $n = 6, 7, 8, \dots, 39$. The results are given in table 5.

With the possible exception of the sequence $\{d_2\}$, the sequences for the various amplitudes appear to be converging. Various other values for the exponents were also tried, including a square-root correction-to-scaling term. In all cases the convergence was dramatically worsened by such changes. Indeed, with a square-root correction-to-scaling exponent, a number of sequences appeared to diverge rather than converge. However, we have assumed above that the sub-leading term of the singularity on the negative real axis is analytic. If we allow this singularity to be a square root singularity, so that the last term in (5.2) above becomes $\mu^n(-1)^n d_2 n^{-2}$ then the results converge even faster, as shown in table 6.

Table 6. Sequences of amplitude estimates, with the exponent associated with d_2 changed from $-\frac{5}{2}$ to -2 . Refer to equation (5.2).

n	d_2	d_1	b_1	a_2	a_1
29	0.0246	-0.1902	-0.2004	0.5585	1.176 99
30	0.0252	-0.1903	-0.2017	0.5589	1.176 99
31	0.0266	-0.1906	-0.1985	0.5580	1.177 00
32	0.0270	-0.1906	-0.1994	0.5583	1.177 00
33	0.0281	-0.1908	-0.1969	0.5576	1.177 00
34	0.0283	-0.1909	-0.1974	0.5577	1.177 00
35	0.0292	-0.1910	-0.1952	0.5571	1.177 01
36	0.0293	-0.1910	-0.1955	0.5572	1.177 01
37	0.0301	-0.1912	-0.1937	0.5568	1.177 01
38	0.0301	-0.1912	-0.1939	0.5568	1.177 01
39	0.0308	-0.1912	-0.1923	0.5564	1.177 02

From these tables we estimate $a_1 \approx 1.1771$, $a_2 \approx 0.554$, $b_1 \approx -0.19$, $d_1 \approx -0.19$, where errors are expected to be confined to the last quoted digit in each case. Repeating the above calculations with a critical point shifted by twice the confidence limit quoted does not change these amplitude estimates.

This then completes our numerical study of the generating function for self-avoiding walks.

Acknowledgments

One of us (ARC) would like to thank the A O Capell, Wyselaskie and Daniel Curdie scholarships. We would like to acknowledge the support of ACCI and The University of Melbourne for the provision of the computers on which these calculations were performed. Mr Glenn Wightwick of IBM (Australia) provided valuable help in running large jobs on the IBM 3090, and we wish to express our thanks for this assistance. This work was supported by the Australian Research Council.

References

- Enting I G 1980 *J. Phys. A: Math. Gen.* **13** 3713
Enting I G and Guttmann A J 1980 *J. Phys. A: Math. Gen.* **13** 1043
— 1985 *J. Phys. A: Math. Gen.* **18** 1007
— 1989 *J. Phys. A: Math. Gen.* **22** 1371
— 1990 *J. Stat. Phys.* **58** 475
— 1992 *J. Phys. A: Math. Gen.* **25** 2791
Guttmann A J 1987 *J. Phys. A: Math. Gen.* **20** 1839
— 1989 *Phase Transitions and Critical Phenomena* vol 13, ed C Domb and J Lebowitz (New York: Academic)
Guttmann A J and Enting I G 1988 *J. Phys. A: Math. Gen.* **21** L165
Guttmann A J and Wang J S 1991 *J. Phys. A: Math. Gen.* **24** 3107
Guttmann A J and Whittington S G 1978 *J. Phys. A: Math. Gen.* **11** 721
MacDonald D, Hunter D L, Kelly K and Jan N 1992 *J. Phys. A: Math. Gen.* **25** 1429
Masand B, Wilensky U, Maffar J P and Redner S 1992 *J. Phys. A: Math. Gen.* **25** L365
Nienhuis B 1982 *Phys. Rev. Lett.* **49** 1062
Nienhuis B 1984 *J. Stat. Phys.* **34** 731
Sykes M F, Guttmann A J, Roberts P D and Watts M G 1972 *J. Phys. A: Math. Gen.* **5** 653